

On the sensitivity of the \mathcal{H}_∞ norm of systems described by delay differential algebraic equations

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Abstract: We consider delay differential algebraic equations (DDAEs) to model interconnected systems with time-delays. The DDAE framework does not require any elimination techniques and can directly deal with any interconnection of systems and controllers with time-delays. In this framework, we analyze the properties of the \mathcal{H}_∞ norm of systems described by delay differential algebraic equations. We show that the standard \mathcal{H}_∞ norm may be sensitive to arbitrarily small delay perturbations. We introduce the strong \mathcal{H}_∞ norm which is insensitive to small delay perturbations and describe its properties. We conclude that the strong \mathcal{H}_∞ norm is more appropriate in any practical control application compared to the standard \mathcal{H}_∞ norm for systems with time-delays whenever there are high-frequency paths in control loops.

Keywords: h-infinity norm, strong h-infinity norm, computational methods, time-delay, interconnected systems, delay differential algebraic equations.

1. INTRODUCTION

In robust control applications, the design requirements are usually defined in terms of \mathcal{H}_∞ norms of the closed-loop functions including the plant, the controller and weights for uncertainties and disturbances Zhou et al. (1995). The properties and robust computational methods of the \mathcal{H}_∞ norm of closed-loop functions are essential in a computer aided control system design. The properties of \mathcal{H}_∞ norm for finite dimensional multi-input-multi-output systems are well-known and reliable numerical methods for the \mathcal{H}_∞ norm computation are available Boyd and Balakrishnan (1990); Bruinsma and Steinbuch (1990).

We analyze the sensitivity of the \mathcal{H}_∞ norm of systems described by delay differential algebraic equations. An important motivation for systems under consideration stems from the fact that interconnected systems with delays can be naturally modeled by state-space representation of the form

$$\begin{cases} E\dot{x}(t) = A_0x(t) + \sum_{i=1}^m A_ix(t - \tau_i) + Bw(t), \\ z = Cx(t). \end{cases} \quad (1)$$

The time-delays τ_i , $i = 1, \dots, m$ are positive real numbers. The system matrices are E and A_i , $i = 0, \dots, m$ are real-valued square matrices and other system matrices with the capital letters are real-valued matrices with appropriate dimensions. The input w and output z are disturbances and signals to be minimized to achieve design requirements and some of system matrices may include the controller parameters.

The system with the closed-loop equations (1) represents all interesting cases of the feedback connection of a time-

delay plant and a controller. The transformation of the closed-loop system to this form can be easily done by first augmenting the system equations of the plant and controller. As we shall see, this augmented system can subsequently be brought in the form (1) by introducing slack variables to eliminate input/output delays and direct feedthrough terms in the closed-loop equations. Hence, the resulting system of the form (1) is obtained directly without complicated elimination techniques, that may even not be possible in the presence of time-delays. It can serve as a standard form for the development of control design and software.

By interconnecting systems and controller high frequency paths could be created in control loops, which may lead to sensitivity problems with respect to the delays and delay perturbations. Therefore it is important to take the sensitivity explicitly into account in the design. We will illustrate that the \mathcal{H}_∞ norm of the transfer function from w to z in (1) may be sensitive to arbitrarily small delay changes. Since small modeling errors are inevitable in any practical design we are interested in the smallest upper bound of the \mathcal{H}_∞ norm that is insensitive to small delay changes. Inspired by the concept of strong stability of neutral equations Hale and Verduyn Lunel (2002), this leads us to the introduction of the concept of *strong \mathcal{H}_∞ norms* for DDAEs. Several properties of the strong \mathcal{H}_∞ norm are shown and a computational formula is obtained. The theory derived can be considered as the dual of the theory of strong stability as elaborated in Hale and Verduyn Lunel (2002); Michiels et al. (2002); Michiels and Vyhlídal (2005); Michiels et al. (2009) and the references therein.

Assumption 5. The matrix $U^T A_0 V$ is nonsingular.

In order to motivate Assumption 5, we note that the equations (1) can be separated into coupled delay differential and delay difference equations. When we define

$$\mathbf{U} = [U^\perp U], \quad \mathbf{V} = [V^\perp V],$$

a pre-multiplication of (1) with \mathbf{U}^T and the substitution

$$x = \mathbf{V} [x_1^T x_2^T]^T,$$

with $x_1(t) \in \mathbb{R}^{n-m}$ and $x_2(t) \in \mathbb{R}^m$, yield the coupled equations

$$\begin{cases} E^{(11)} \dot{x}_1(t) = \sum_{i=0}^m A_i^{(11)} x_1(t - \tau_i) + \sum_{i=0}^m A_i^{(12)} x_2(t - \tau_i) + B_1 w(t), \\ 0 = A_0^{(22)} x_2(t) + \sum_{i=1}^m A_i^{(22)} x_2(t - \tau_i) + \sum_{i=0}^m A_i^{(21)} x_1(t - \tau_i) + B_2 w(t), \\ y(t) = C_1 x_1(t) + C_2 x_2(t), \end{cases} \quad (5)$$

where

$$\begin{aligned} A_i^{(11)} &= U^{\perp T} A_i V^\perp, \quad A_i^{(12)} = U^{\perp T} A_i V, \\ A_i^{(21)} &= U^T A_i V^\perp, \quad A_i^{(22)} = U^T A_i V, \quad i = 0, \dots, m, \end{aligned} \quad (6)$$

and

$$\begin{aligned} E^{(11)} &= U^{\perp T} E V^\perp, \quad B_1 = U^{\perp T} B, \quad B_2 = U^T B, \\ C_1 &= C V^\perp, \quad C_2 = C V. \end{aligned} \quad (7)$$

Matrix $E^{(11)}$ in (5) is invertible, following from

$$n - m = \text{rank}(E) = \text{rank}(\mathbf{U}^T E \mathbf{V}) = \text{rank}(E^{(11)}).$$

In addition, matrix $A_0^{(22)}$ is invertible, following from Assumption 5.

The equations (5) with $w \equiv 0$ are semi-explicit delay differential algebraic equations of index 1, because delay differential equations are obtained by differentiating the second equation. This precludes the occurrence of impulsive solutions Fridman and Shaked (2002). Moreover, the invertibility of $A_0^{(22)}$ prevents that the equations are of *advanced* type and, hence, non-causal. This further motivates why Assumption 5 is natural in the delay case considered, although it restricts the index to one (for a general treatment in the delay free case, see for instance Stykel (2002) and the references therein).

We also make the following assumption.

Assumption 6. The zero solution of system (1), with $w \equiv 0$, is strongly exponentially stable.

Strong exponential stability refers to the fact that the asymptotic stability of the null solution is robust against small delay perturbations Hale and Verduyn Lunel (2002); Michiels et al. (2009). Due to the modeling errors and uncertainty, the delays of the time-delay model are typically not exactly known and this type of stability is required in practice. The stability of the closed-loop system (1) is a necessary assumption since the \mathcal{H}_∞ norm is defined for stable systems only.

Transfer functions

From (5) we can write the transfer function of the system (1) as

$$T(\lambda) := C \left(\lambda E - A_0 - \sum_{i=1}^m A_i e^{-\lambda \tau_i} \right)^{-1} B, \quad (8)$$

$$= [C_1 \ C_2] \begin{bmatrix} \lambda E^{(11)} - A_{11}(\lambda) & -A_{12}(\lambda) \\ -A_{21}(\lambda) & -A_{22}(\lambda) \end{bmatrix}^{-1} \begin{bmatrix} B_1 \\ B_2 \end{bmatrix}, \quad (9)$$

with

$$A_{kl}(\lambda) = \sum_{i=0}^m A_i^{(kl)} e^{-\lambda \tau_i}, \quad k, l \in \{1, 2\}.$$

We define the *asymptotic* transfer function of the system (1) as

$$\begin{aligned} T_a(\lambda) &:= -C V \left(U^T A_0 V + \sum_{i=1}^m U^T A_i V e^{-\lambda \tau_i} \right)^{-1} U^T B \\ &= -C_2 A_{22}(\lambda)^{-1} B_2. \end{aligned} \quad (10)$$

The terminology stems from the fact that the transfer function T and the asymptotic transfer function T_a converge to each other for high frequencies, as precized in the following Proposition.

Proposition 7. $\forall \gamma > 0, \exists \Omega > 0: \sigma_1(T(j\omega) - T_a(j\omega)) < \gamma, \forall \omega > \Omega.$

Proof. The assertion follows from the explicit expression for the inverse of the two-by-two block matrix in (9), combined with the property that

$$\sup_{\Re(\lambda) \geq 0} \left\| (A_{22}(\lambda))^{-1} \right\|_2 \quad (12)$$

if finite. The latter is due to Assumption 6. \square

The \mathcal{H}_∞ norm of the transfer function T of the *stable* system (1), is defined as

$$\|T(j\omega)\|_\infty := \sup_{\omega \in j\mathbb{R}} \sigma_1(T(j\omega)). \quad (13)$$

Similarly we can define \mathcal{H}_∞ norm of T_a .

4. STRONG \mathcal{H}_∞ NORM OF TIME-DELAY SYSTEMS

We now analyze continuity properties of the \mathcal{H}_∞ norm of the transfer function T with respect to the delay parameters. The function

$$\boldsymbol{\tau} \in (\mathbb{R}_0^+)^m \mapsto \|T(j\omega, \boldsymbol{\tau})\|_\infty \quad (14)$$

is, in general, not continuous, which is inherited from the behavior of the asymptotic transfer function, T_a , more precisely the function

$$\boldsymbol{\tau} \in (\mathbb{R}_0^+)^m \mapsto \|T_a(j\omega, \boldsymbol{\tau})\|_\infty. \quad (15)$$

We start with a motivating example

Example 8. Let the transfer function T be defined as

$$T(\lambda) = \frac{\lambda + 2}{\lambda(1 - 0.25e^{-\lambda \tau_1} + 0.5e^{-\lambda \tau_2}) + 1} \quad (16)$$

where $(\tau_1, \tau_2) = (1, 2)$. The transfer function T is stable, its \mathcal{H}_∞ norm is 2.6422 achieved at $\omega = 1.6598$ and the maximum singular value plot is given in Figure 1 (on the left). The high frequency behavior is described by the asymptotic transfer function

$$T_a(\lambda) = \frac{1}{(1 - 0.25e^{-\lambda \tau_1} + 0.5e^{-\lambda \tau_2})}, \quad (17)$$

whose \mathcal{H}_∞ norm is equal to 2.0320, which is less than $\|T(j\omega)\|_\infty$. However, when the first time delay is perturbed

to $\tau_1 = 0.99$, the \mathcal{H}_∞ norm of the transfer function T is 3.9993, reached at $\omega = 158.6578$, see Figure 1 (on the right). The \mathcal{H}_∞ norm of T is quite different from that for $(\tau_1, \tau_2) = (1, 2)$. A closer look at the maximum singular value plot of the asymptotic transfer function T_a in Figure 2 (on the left) shows that the sensitivity is due to the transfer function T_a .

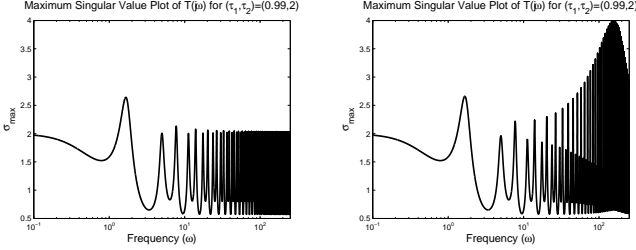


Fig. 1. The maximum singular value plot of $T(j\omega)$ for $(\tau_1, \tau_2) = (1, 2)$ (left) and $(\tau_1, \tau_2) = (0.99, 2)$ (right) as a function of ω .

Even if the first delay is perturbed slightly to $\tau_1 = 0.999$, the problem is not resolved, indicating that the functions (14) and (15) are discontinuous at $(\tau_1, \tau_2) = (1, 2)$. The \mathcal{H}_∞ norm of the transfer function T for $(\tau_1, \tau_2) = (0.999, 2)$ is namely given by 3.9998, and the peak value is reached at $\omega = 1566.0816$. The corresponding asymptotic transfer function T_a is shown in Figure 2 (on the right). When the delay perturbation tends to zero, the frequency where the maximum in the singular value plot of the asymptotic transfer function T_a is achieved moves towards infinity.

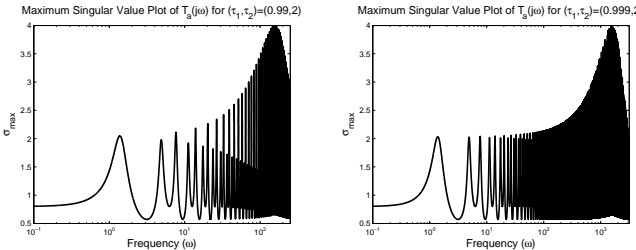


Fig. 2. The maximum singular value plot of $T_a(j\omega)$ for $(\tau_1, \tau_2) = (0.99, 2)$ (left) and $(\tau_1, \tau_2) = (0.999, 2)$ (right) as a function of ω .

The above example illustrates that the \mathcal{H}_∞ norm of the transfer function T may be sensitive to *infinitesimal* delay changes. Since this property is related to the behavior of the transfer function at high frequencies and, hence, the asymptotic transfer function T_a , we first study the properties of the function (15).

Since small modeling errors and uncertainty are inevitable in a practical design, we wish to characterize the smallest upper bound for the \mathcal{H}_∞ norm of the asymptotic transfer function T_a which is *insensitive* to small delay changes.

Definition 9. For $\tau \in (\mathbb{R}_0^+)^m$, let the strong \mathcal{H}_∞ norm of T_a , $\|T_a(j\omega, \tau)\|_\infty$, be defined as

$$\|T_a(j\omega, \tau)\|_\infty := \lim_{\epsilon \rightarrow 0^+} \sup \{ \|T_a(j\omega, \tau_\epsilon)\|_\infty : \tau_\epsilon \in \mathcal{B}(\tau, \epsilon) \cap (\mathbb{R}^+)^m \}.$$

Several properties of this upper bound on $\|T_a(j\omega, \tau)\|_\infty$ are listed below.

Proposition 10. The following assertions hold:

(1) for every $\tau \in (\mathbb{R}_0^+)^m$, we have

$$\|T_a(j\omega, \tau)\|_\infty = \max_{\theta \in [0, 2\pi]^m} \sigma_1(\mathbb{T}_a(\theta)), \quad (18)$$

where

$$\begin{aligned} \mathbb{T}_a(\theta) &= C_2 \left(-A_0^{(22)} - \sum_{i=1}^m A_i^{(22)} e^{-j\theta_i} \right)^{-1} B_2, \quad (19) \\ &= CV \left(-U^T A_0 V - \sum_{i=1}^m U^T A_i V e^{-j\theta_i} \right)^{-1} U^T B. \end{aligned}$$

(2) $\|T_a(j\omega, \tau)\|_\infty \geq \|T_a(j\omega, \tau)\|_\infty$ for all delays τ ;

(3) $\|T_a(j\omega, \tau)\|_\infty = \|T_a(j\omega, \tau)\|_\infty$ for rationally independent¹ τ .

Proof. We always have

$$(e^{-j\omega\tau_1}, \dots, e^{-j\omega\tau_m}) \in \{(e^{-j\theta_1}, \dots, e^{-j\theta_m}) : \theta_i \in [0, 2\pi], i = 1, \dots, m\},$$

implying

$$\|T(j\omega, \tau)\|_\infty \leq \max_{\theta \in [0, 2\pi]^m} \sigma_1(\mathbb{T}_a(\theta)). \quad (20)$$

For any $\epsilon > 0$ in Definition 9, there exists $\tau_\epsilon = [\tau_{\epsilon,1}, \dots, \tau_{\epsilon,m}]$ rationally independent in $\mathcal{B}(\tau, \epsilon) \cap (\mathbb{R}^+)^m$. By Theorem 2.1 in Michiels et al. (2002), given rationally independent time delays τ_ϵ and for $\theta = [\theta_1, \dots, \theta_m]$ arbitrary, there exists a sequence of real numbers $\{\omega_n\}_{n \geq 1}$ such that

$$\lim_{n \rightarrow \infty} \max_{1 \leq i \leq m} |e^{-j\omega_n \tau_{\epsilon,i}} - e^{-j\theta_i}| = 0.$$

It follows that

$$\begin{aligned} \text{closure}\{(e^{-j\omega\tau_{\epsilon,1}}, \dots, e^{-j\omega\tau_{\epsilon,m}}) : \omega \in \mathbb{R}\} = \\ \{(e^{-j\theta_1}, \dots, e^{-j\theta_m}) : \theta_i \in [0, 2\pi], i = 1, \dots, m\}, \end{aligned}$$

implying

$$\|T(j\omega, \tau_\epsilon)\|_\infty = \max_{\theta \in [0, 2\pi]^m} \sigma_1(\mathbb{T}_a(\theta)). \quad (21)$$

The assertions follow from (20) and (21). \square

Formula (18) in Proposition 10 shows that the strong \mathcal{H}_∞ norm is independent of the delay values. The formula further naturally leads to a computational scheme based on sweeping on θ intervals. This approximation can be corrected by solving a set of nonlinear equations. Numerical computation details are presented in Gumussoy and Michiels (2010).

We now come back to the properties of the transfer function (14) of the system (1). As we have illustrated, a discontinuity of the function (15) may carry over to the function (14). Therefore, we define the strong \mathcal{H}_∞ norm of the transfer function T in a similar way.

Definition 11. For $\tau \in (\mathbb{R}_0^+)^m$, the strong \mathcal{H}_∞ norm of T , $\|T(j\omega, \tau)\|_\infty$, is given by

$$\|T(j\omega, \tau)\|_\infty := \lim_{\epsilon \rightarrow 0^+} \sup \{ \|T(j\omega, \tau_\epsilon)\|_\infty : \tau_\epsilon \in \mathcal{B}(\tau, \epsilon) \cap (\mathbb{R}^+)^m \}. \quad (22)$$

¹ The m components of $\tau = (\tau_1, \dots, \tau_m)$ are rationally independent if and only if $\sum_{k=1}^m z_k \tau_k = 0$, $z_k \in \mathbb{Z}$ implies $z_k = 0$, $\forall k = 1, \dots, m$. For instance, two delays τ_1 and τ_2 are rationally independent if their ratio is an irrational number.

The following main theorem describes, among others, the desirable property that, in contrast to the \mathcal{H}_∞ norm, the strong \mathcal{H}_∞ norm *continuously* depends on the delay parameters. The proof makes use of the technical results in Section 7 of the appendix.

Theorem 12. The strong \mathcal{H}_∞ norm of the delay differential algebraic system (1) satisfies

$$\| \|T(j\omega, \boldsymbol{\tau})\| \|_\infty = \max(\|T(j\omega, \boldsymbol{\tau})\|_\infty, \| \|T_a(j\omega, \boldsymbol{\tau})\| \|_\infty), \quad (23)$$

where T and T_a are the transfer function (8) and the asymptotic transfer function (10). In addition, the function

$$\boldsymbol{\tau} \in (\mathbb{R}_0^+)^m \mapsto \| \|T(j\omega, \boldsymbol{\tau})\| \|_\infty \quad (24)$$

is continuous.

Proof. Lemma 14 implies that the function (14) is continuous at delay values where

$$\|T(j\omega, \boldsymbol{\tau})\|_\infty > \| \|T_a(j\omega, \boldsymbol{\tau})\| \|_\infty. \quad (25)$$

This property, along with the fact that $\| \|T_a(j\omega, \boldsymbol{\tau})\| \|_\infty$ is independent of $\boldsymbol{\tau}$ (see Proposition 10), lead to the assertion (23) and the continuity of (24) under the condition (25). In the other case the assertions follow from Lemma 15. \square

The explicit expression (23) lays at the basis of an algorithm to compute the strong \mathcal{H}_∞ norm presented in the accompanying paper Gumussoy and Michiels (2010).

5. NUMERICAL EXAMPLE

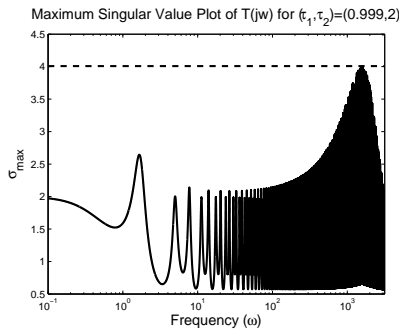


Fig. 3. The maximum singular value plot of T (16): $\|T(j\omega, \boldsymbol{\tau})\|_\infty < \| \|T_a(j\omega, \boldsymbol{\tau})\| \|_\infty$ case.

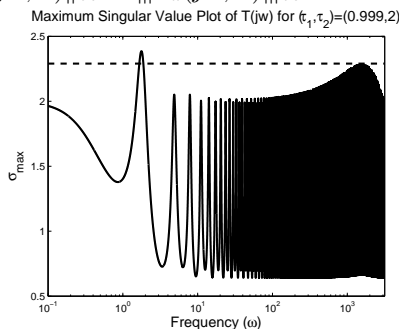


Fig. 4. The maximum singular value plot of T (26): $\|T(j\omega, \boldsymbol{\tau})\|_\infty > \| \|T_a(j\omega, \boldsymbol{\tau})\| \|_\infty$ case.

By (23), the strong \mathcal{H}_∞ norm of the transfer function T is determined by either the \mathcal{H}_∞ norm of T or the strong \mathcal{H}_∞ norm of T_a . We illustrate both cases.

Given the transfer function T (16), the strong \mathcal{H}_∞ norm of its asymptotic transfer function T_a is equal to 4 (indicated

as a dashed line) and the \mathcal{H}_∞ norm of T is 2.6422 as shown in Figure 3. Then the strong \mathcal{H}_∞ norm of T (16) is equal to the strong \mathcal{H}_∞ norm of (17), namely 4.

As a second example, consider the transfer function

$$T(\lambda, \boldsymbol{\tau}) := \frac{\lambda + 2}{\lambda(1 - 1/16e^{-\lambda\tau_1} + 1/2e^{-\lambda\tau_2}) + 1}, \quad (26)$$

with $\boldsymbol{\tau} = (1, 2)$, and its asymptotic transfer function

$$T_a(\lambda, \boldsymbol{\tau}) := \frac{1}{(1 - 1/16e^{-\lambda\tau_1} + 1/2e^{-\lambda\tau_2})}. \quad (27)$$

Figure 4 shows that the strong \mathcal{H}_∞ norm of T (26) is equal to the \mathcal{H}_∞ norm of T (26). Note that the strong \mathcal{H}_∞ norm of the asymptotic transfer function can be used as the first level to compute the strong \mathcal{H}_∞ norm in well-known level set methods Boyd and Balakrishnan (1990); Bruinsma and Steinbuch (1990).

6. CONCLUDING REMARKS

We analyzed the sensitivity of the \mathcal{H}_∞ norm of interconnected systems with time-delays. We showed that a very broad class of interconnected retarded and/or neutral systems can be brought in the standard form (1) in a systematic way. Input/output delays and direct feedthrough terms can be dealt with by introducing slack variables. An additional advantage in the context of control design is the linearity of the closed loop matrices w.r.t. the controller parameters.

We showed the sensitivity of the \mathcal{H}_∞ norm w.r.t. small delay perturbations and introduced the *strong \mathcal{H}_∞ norm* for DDAEs inline with the notion of strong stability. We analyzed its continuity properties derived as an explicit expression. The given properties are illustrated on numerical examples.

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7. SOME TECHNICAL LEMMAS

Lemma 13. For all $\gamma > 0$, there exist numbers $\epsilon > 0$ and $\Omega > 0$ such that

$$\sigma_1(T(j\omega, \mathbf{r}) - T_a(j\omega, \mathbf{r})) < \gamma$$

for all $\omega > \Omega$ and $\mathbf{r} \in \mathcal{B}(\boldsymbol{\tau}, \epsilon) \cap (\mathbb{R}^+)^m$.

Proof. The uniformity of the bound γ w.r.t. small delay perturbations stems from the fact that the bound (12) is a continuous function of the delays $\boldsymbol{\tau}$ at their nominal values. The latter is implied by the *strong* stability assumption (Assumption 6). \square

Lemma 14. Let $\xi > \|T_a(j\omega, \boldsymbol{\tau})\|_\infty$ hold. Then there exist real numbers $\epsilon > 0$, $\Omega > 0$ and an integer N such that for any $\mathbf{r} \in \mathcal{B}(\boldsymbol{\tau}, \epsilon) \cap (\mathbb{R}^+)^m$, the number of frequencies $\omega^{(i)}$ such that

$$\sigma_k(T(j\omega^{(i)}, \mathbf{r})) = \xi, \quad (28)$$

for some $k \in \{1, \dots, n\}$, is smaller than N , and, moreover, $|\omega^{(i)}| < \Omega$.

Proof. For any (fixed) value of $\xi > 0$ and delays \mathbf{r} , the relation

$$\sigma_k(T(j\omega), \mathbf{r}) = \xi \quad (29)$$

holds for some $\omega \in \mathbb{R}$ and $k \in \{1, \dots, n\}$ if and only if $\lambda = j\omega$ is a zero of the function

$$\det \left(\begin{bmatrix} \lambda E - A_0 - \sum_{i=1}^m A_i e^{-\lambda r_i} & -\frac{1}{\xi} B B^T \\ \frac{1}{\xi} C C^T & \lambda E^T + A_0^T + \sum_{i=1}^m A_i^T e^{\lambda r_i} \end{bmatrix} \right). \quad (30)$$

This result is a variant of Lemma 2.1 of Michiels and Gumussoy (2010) to which we refer for the proof.

Now take $\xi > \|T_a(j\omega, \boldsymbol{\tau})\|_\infty$. From Lemma 13, and taking into account that $\|T_a(j\omega, \boldsymbol{\tau})\|_\infty$ does not depend on $\boldsymbol{\tau}$ (see Proposition 10) it follows that there exists numbers $\epsilon > 0$ and $\Omega > 0$ such that all ω satisfying (29) for some $k \in \{1, \dots, n\}$ and $\mathbf{r} \in \mathcal{B}(\boldsymbol{\tau}, \epsilon) \cap (\mathbb{R}^+)^m$ also satisfy $|\omega| < \Omega$. This proves one statement. At the same time $\lambda = j\omega$ must be a zero of the analytic function (30). The other statement is due to the fact that an analytic function only has finitely many zeros in a compact set. \square

Lemma 15. The following implication holds

$$\|T(j\omega, \boldsymbol{\tau})\|_\infty \leq \|T_a(j\omega, \boldsymbol{\tau})\|_\infty \Rightarrow \|T(j\omega, \boldsymbol{\tau})\|_\infty = \|T_a(j\omega, \boldsymbol{\tau})\|_\infty.$$

Proof. For every $\epsilon > 0$ there exist delays $\boldsymbol{\tau}_0$ and a frequency ω_0 such that

$$\|\boldsymbol{\tau}_0 - \boldsymbol{\tau}\| < \epsilon/2, \quad \sigma_1(T_a(j\omega_0, \boldsymbol{\tau}_0)) \geq \|T_a(j\omega, \boldsymbol{\tau})\|_\infty - \epsilon/2.$$

In addition, there exist commensurate delays

$$\boldsymbol{\tau}_r = (n_1/s, \dots, n_m/s), \quad (31)$$

with $(n_1, \dots, n_m, s) \in \mathbb{N}^{m+1}$ such that

$$\|\boldsymbol{\tau}_r - \boldsymbol{\tau}_0\| < \epsilon/2, \\ |\sigma_1(T_a(j\omega_0, \boldsymbol{\tau}_r)) - \sigma_1(T_a(j\omega_0, \boldsymbol{\tau}_0))| \leq \epsilon/2.$$

Thus, for all $\epsilon > 0$ there exist commensurate delays (31) and a frequency ω_0 satisfying

$$\|\boldsymbol{\tau}_r - \boldsymbol{\tau}\| < \epsilon, \quad \sigma_1(T_a(j\omega_0, \boldsymbol{\tau}_r)) \geq \|T_a(j\omega, \boldsymbol{\tau})\|_\infty - \epsilon.$$

From the fact that

$$T_a(j\omega_0, \boldsymbol{\tau}_r) = T_a(j(\omega_0 + 2\pi sk), \boldsymbol{\tau}_r)$$

for all $k \geq 1$ and Lemma 13, we conclude that

$$\|T(j\omega, \boldsymbol{\tau})\|_\infty \geq \|T_a(j\omega, \boldsymbol{\tau})\|_\infty. \quad (32)$$

Now take a level $\xi > \|T_a(j\omega, \boldsymbol{\tau})\|_\infty$, and let ϵ and Ω be determined by the assertion of Lemma 14. From the assumption $\|T(j\omega, \boldsymbol{\tau})\|_\infty \leq \|T_a(j\omega, \boldsymbol{\tau})\|_\infty$ and the relation between (29) and (30) it follows that the function (30) has no zeros on the imaginary axis for $\mathbf{r} = \boldsymbol{\tau}$. Because the function (30) is analytic and all potential imaginary axis zeros have modulus smaller than Ω whenever $\mathbf{r} \in \mathcal{B}(\boldsymbol{\tau}, \epsilon) \cap (\mathbb{R}^+)^m$, we conclude that there exists a number $\epsilon_2 > 0$ such that the function (30) has no imaginary axis eigenvalues whenever $\mathbf{r} \in \mathcal{B}(\boldsymbol{\tau}, \epsilon_2) \cap (\mathbb{R}^+)^m$. Equivalently, $T(j\omega, \mathbf{r})$ has no singular values equal to ξ whenever $\mathbf{r} \in \mathcal{B}(\boldsymbol{\tau}, \epsilon_2) \cap (\mathbb{R}^+)^m$. This proves that the left and the right hand side of (32) are equal. \square